

Integrated Q-switched lasing element in the NIR with transition metal dichalcogenide gain and graphene saturable absorption

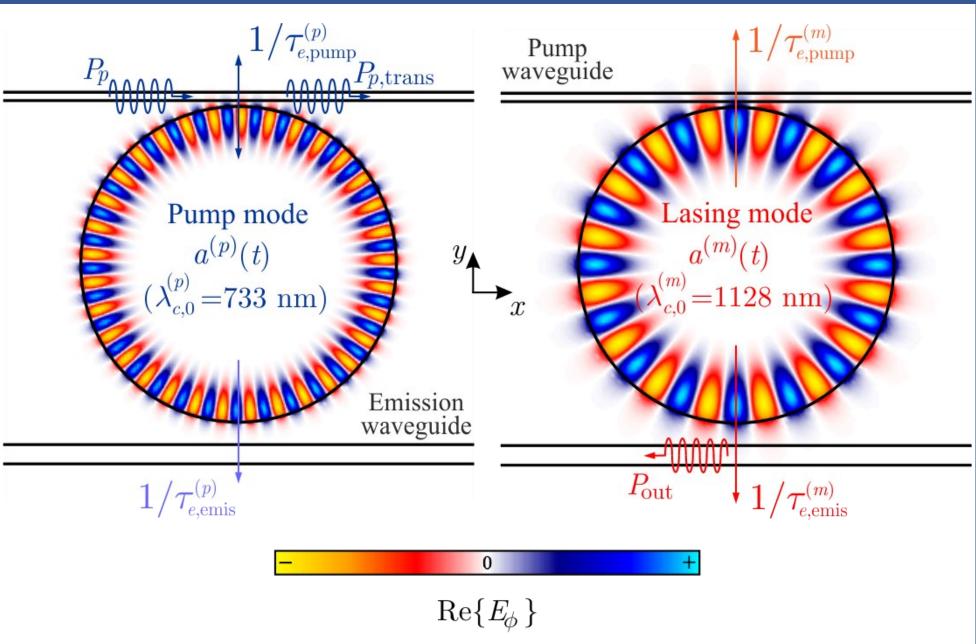
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I. Introduction

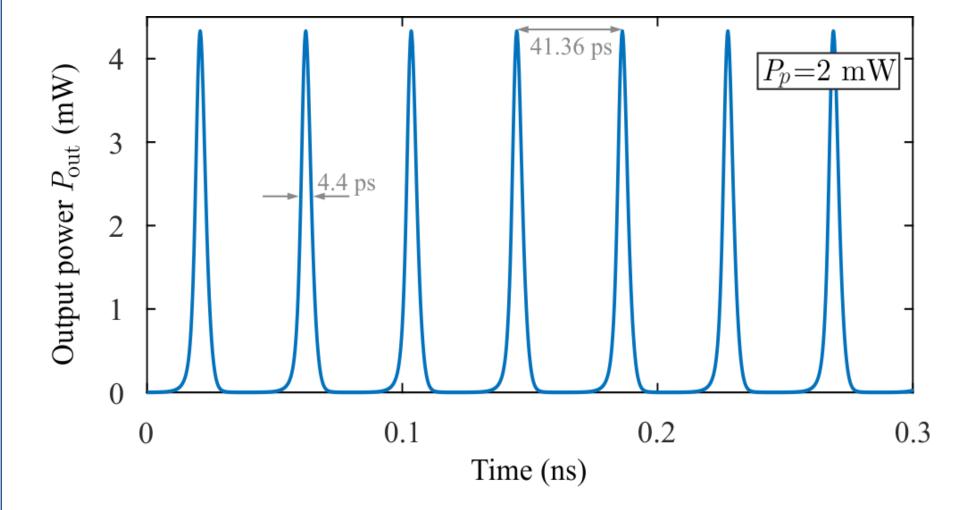
- Motivation:
 - Bring passive Q-switching at the nanoscale aiming to develop efficient integrated silicon-on-insulator (SOI)compatible pulsed laser sources in the near infrared (NIR).
 - Exploit the novel luminescence, linear and nonlinear properties of contemporary two-dimensional (2D) materials to provide *both* the gain and saturable absorption (SA) mechanisms.
- Objective:
 - Propose, design, and numerically analyze an integrated nanophotonic passively Q-switched lasing element based on

III. Integrated laser structure



V. Laser pulsed operation

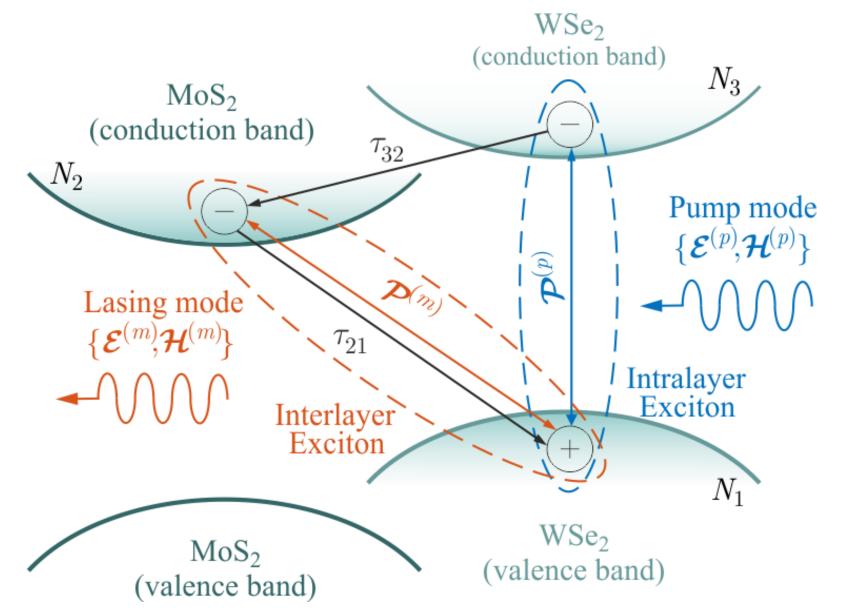
- Due to the fulfillment of the critical coupling condition for the pump mode, the CW lasing threshold is just 25 µW.
- Q-switched operation with infinite extinction ratio is obtained from pump power levels just above the lasing threshold up to 3.5 mW.



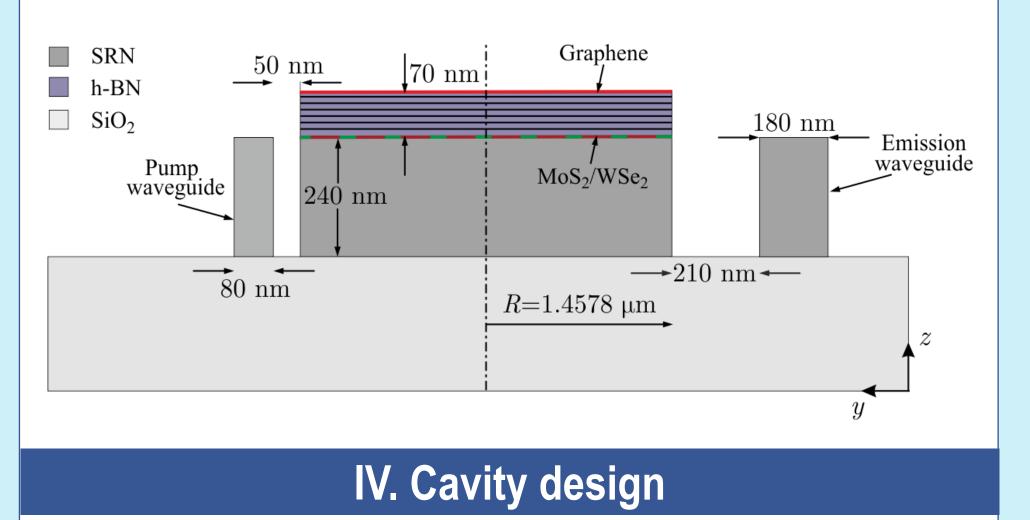
a **silicon-rich nitride (SRN) disk resonator**, enhanced with gain provided by a **MoS₂/WSe₂** transition metal dichalcogenide (TMD) **hetero-bilayer** and SA by a **graphene monolayer**.

II. Computational framework

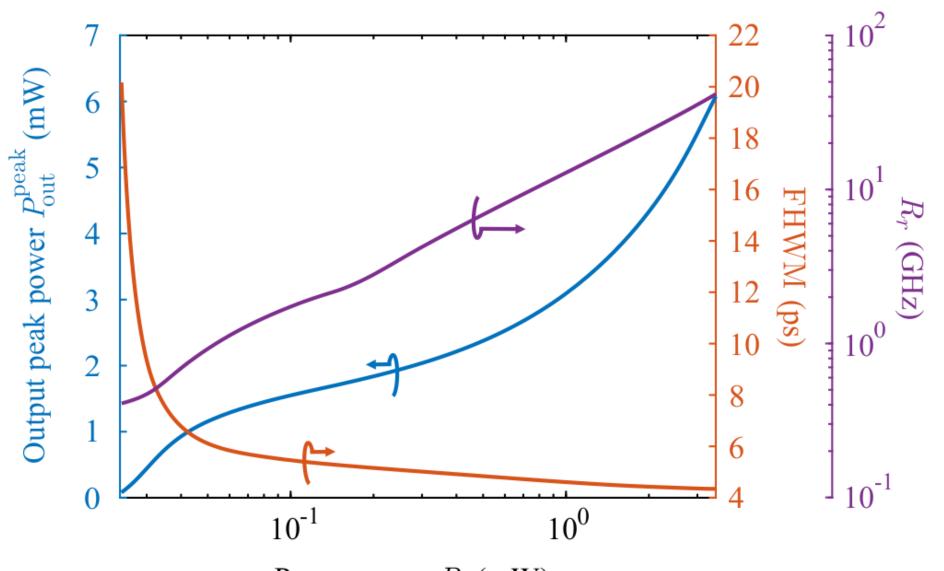
- Accurate and efficient temporal coupled-mode theory (CMT) framework, rigorously extracted from the Maxwell-Bloch equations using first-order perturbation theory and utilizing the standard slowly-varying envelope and rotating-wave approximations [1].
 - Capable of treating both **sheet** and **bulk gain media**.
 - Capable of incorporating nonlinear effects including SA in 2D materials [2].



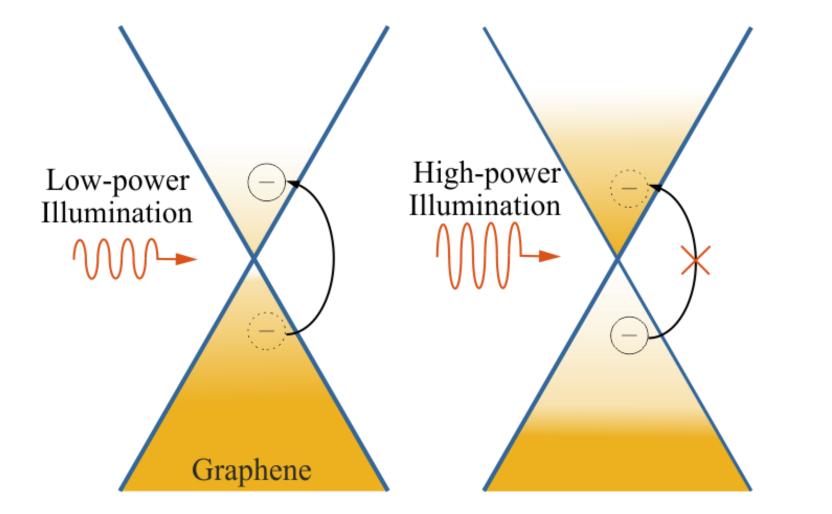
- SRN disk resonator on SiO₂ substrate supporting tightly-confined WGMs with high Q-factors.
 - SRN is **transparent** at *both* the lasing and pump wavelengths.
- MoS₂/WSe₂ and graphene are patterned in disks matching the resonator and placed on top of it, separated by a spacer of (multilayer) hexagonal boron nitride (h-BN).
 - h-BN allows for tailoring the modal overlap with the 2D gain and SA media.
- To efficiently excite the pump mode and extract the emitted light from the lasing mode, the resonator is side-coupled to two dissimilar bus waveguides.
 - \circ Each mode is predominately coupled only to one waveguide.



- The output **peak power** of the Q-switched pulses increases with the pump power from 79 μ W to 6.1 mW. **1 mW peak power** is achieved for just $P_p = 45 \mu$ W.
- **FWHM** is in the **ps** timescale and decreases slowly from 6.35 ps $(P_p = 45 \,\mu\text{W})$ to its lowest value of 4.35 ps $(P_p = 3.5 \,\text{mW})$.
- Repetition rate R_r increases in a practically linear manner with pump power and in a range of two orders of magnitude spanning from 0.41 GHz (2.4 ns) to 41.7 GHz (24 ps).

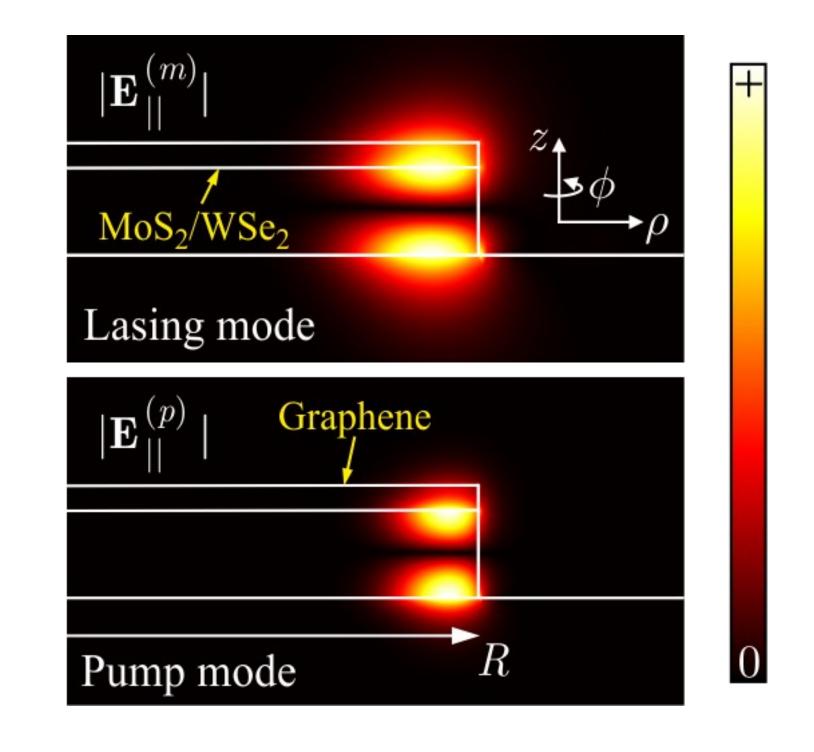


- MoS₂/WSe₂ is a three-level 2D gain medium with long radiative recombination times (~ns).
- Emits light at 1128 nm due to radiative recombination of interlayer excitons after being optically pumped at ~740 nm [3].
- Pumping with guided light to increase absorption efficiency and lower lasing threshold.
- The lasing and pump transitions are described by induced electric polarization fields following the damped Lorentzian oscillator equations and the carrier dynamics by semiclassical carrier-rate equations.



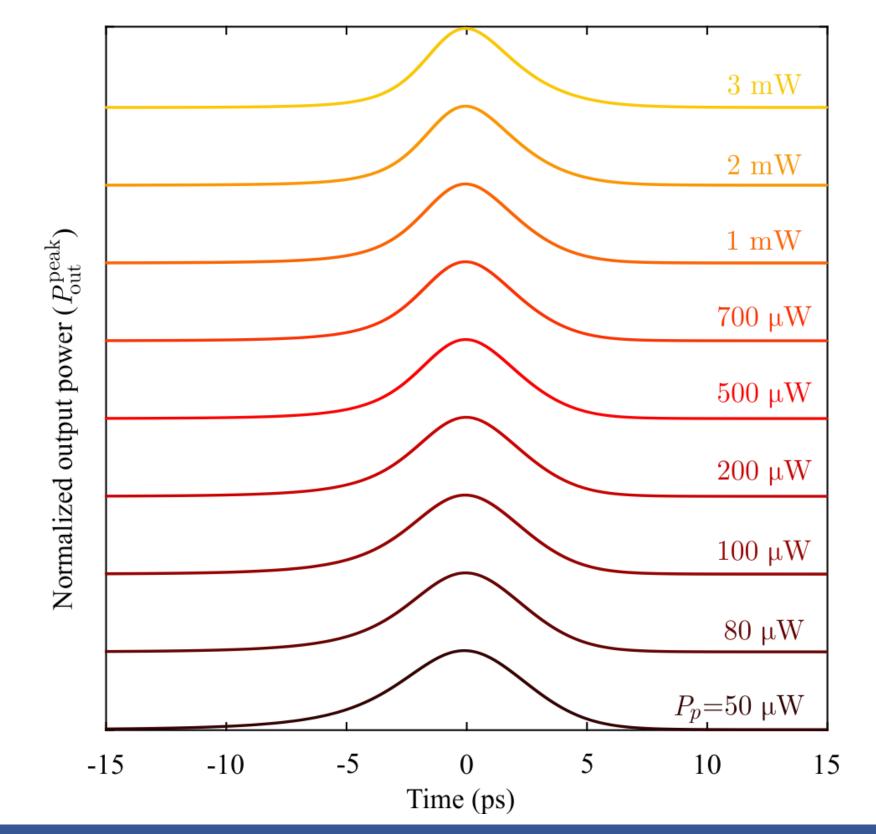
- Graphene possesses ultrafast and low saturation intensity (~1 MW/cm²) SA effect, attributed to Pauli blocking of interband transition [4].
- Instantaneous graphene SA model where only losses are quenched by the lasing mode:
 - σ_0

- The heights of the SRN and h-BN layers in the disk resonator are designed to achieve:
 - 1. Strong light-matter interaction between the lasing mode and both MoS₂/WSe₂ and graphene.
 - 2. Strong light-matter interaction between the pump mode and MoS₂/WSe₂.
- 3. Weak light-matter interaction between the pump mode and graphene.
- For the coupled cavity, we opt for:
 - 1. Over-critical coupling condition between the lasing mode and the emission waveguide to increase the output emitted power.
 - 2. Critical coupling condition between the pump mode and the pump waveguide to maximize absorption efficiency ($\eta_p = 1$), minimizing the lasing threshold.



Pump power P_p (mW)

- The lineshape of the individual Q-switched pulses is highlysymmetric as saturable and non-saturable losses of the lasing mode are equally dominant.
- Transition from pulses with slightly longer leading edges to pulses with longer trailing edges as the pump power increases, revealing the **non-negligible** role of the pump mode in the overall lasing dynamics.



 $\sigma_{\rm gr} = \frac{1}{1 + |\mathbf{E}^{(m)}|^2 / E_{\rm sa^{\dagger}}^2}$

- Overall, the developed CMT framework consists of seven coupled first-order nonlinear ordinary differential equations [two equations for cavity amplitudes (lasing and pump), two equations for induced polarization (lasing and pump), and three carrier rate equations].
 - The coefficients of the CMT equations are evaluated through linear electromagnetic finite element method (FEM) simulations.

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VI. Conclusions

We have **proposed** and thoroughly **studied** a **nanophotonic passively Q-switched lasing element** based on a WGM disk cavity enhanced with **MoS**₂/WSe₂ optically-pumped gain and **graphene** SA. The response of the lasing structure has been accurately **evaluated** through a rigorous **CMT framework**. Overall, the proposed laser has an ultralow lasing threshold and emits pulsed light of mW peak power, ps duration and GHz repetition rate, even for sub-mW pump power levels, rendering it **highly promising** for a variety of **optical communication** and **sensing applications** in the NIR.

VII. References

[1] Nousios *et al.*, Phys. Rev. Appl. **19**, 064027 (2023).
[2] Ataloglou *et al.*, Phys. Rev. A **97**, 063836 (2018).
[3] Liu *et al.*, Sci. Adv. **5**, eaav4506 (2019).
[4] Marini *et al.*, Phys. Rev. B **95**, 125408 (2017).